Detecting topological features in real space and time

Pietro Massignan











The Institute of Photonic

Sciences

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Condensed matter



Plenty of emergent phenomena! But we need to *quantify* these... E.g., how to "detect topology"?



Outline

- Topology in condensed matter
- One-dimensional chiral models
- Quantum Simulation:
 - Topological Anderson Insulator
 - ✦ Interacting fermionic chains





- Quantum Walks ↔ periodically-driven systems:
 - Chiral insulators in 1D QWs
 - Chern insulators in 2D QWs



Topology

 Classification of objects and manifolds under continuous deformations

✓ stretch and bend✗ but don't cut, puncture, or glue

- *Global* properties!
- Genus (# of holes)
- Winding number of a closed path (# of times it encircles a given object)





Hall effect

- Classical Hall effect (1879): when current flows in a 2D material, in presence of an out-of-plane B field, there appears a transverse (Hall) current
- Quantum Hall effect (1980): at low temperatures and high-B, the Hall current is quantized!





- Laughlin (1982): robustness due to topology
- TKNN (1982): Kubo formula links conductivity to *Chern numbers* (topological invariants defined on the occupied bands).

Thouless, Kohmoto, Nightingale & den Nijs Phys. Rev. Lett. (1982)

Hofstadter butterfly

A very simple problem hosting a *fractal* spectrum with topological meaning



Topological insulators

- Insulators in the bulk, with few conducting modes on their edges.
- Protected by the band topology vs. local perturbations, like *disorder* and *defects*.
- Bulk/edge correspondence: edge modes intimately related to topological invariants.

chira

- Enormous progresses in the last 20 years (QSH, 3D TIs., 4D QH, ...)
- Characterization of non-interacting TIs in terms of <u>discrete symmetries</u>
 T: time-reversal
 C: charge-conjugation
 S: chiral
 IQHE, Hofstadter, Chern insulators

 Beyond the periodic table: Mott / crystalline / Anderson / Floquet TIs, ... Winding

 \mathbb{Z}

 \mathbb{Z}_2

0

0

 \mathbb{Z}

 \mathbb{Z}_2

0

 \mathbb{Z}

 \mathbb{Z}_2

 \mathbb{Z}_2

 \mathbb{Z}

 \mathbb{Z}_2

 \mathbb{Z}_2^-

 $2\mathbb{Z}$

0

 \mathbb{Z}

 \mathbb{Z}_2

 \mathbb{Z}_2^2

 $2\mathbb{Z}$

0

0

 \mathbb{Z}

 \mathbb{Z}_2

 $\mathbb{Z}_2^{\tilde{2}}$

0

 $2\mathbb{Z}$

0

0

AI

D

BDI

DIII

AII

CII

С

CI

7

0

 \mathbb{Z}

 \mathbb{Z}_{2}

 \mathbb{Z}_2

0

 $2\mathbb{Z}$

0

0

0

 \mathbb{Z}

0

 $2\mathbb{Z}$

0

0

0

 \mathbb{Z}

 \mathbb{Z}_2

1D chiral systems



polyacetilene [Nobel prize in Chemistry 2000]



ultracold atoms in superlattices [M. Atala *et al.*, Nature Phys. 2013]



[Zeuner *et al.*, PRL 2015]



Cavity polaritons [St. Jean *et al.*, Nature Phot. 2017]



 $t-\Delta$ $t+\Delta$ $t-\Delta$ $t+\Delta$ $t-\Delta$ $t-\Delta$ $t+\Delta$ $t-\Delta$ $t+\Delta$ $t-\Delta$

ultracold atoms in momentum-lattices [Meier *et al.*, Nature Comm. 2016]



SC qubits in mw-cavities [Flurin *et al.*, PRX 2017]

SSH model

• Spinless fermions with staggered tunnelings:

Su, Schrieffer & Heeger Phys. Rev. Lett. (1979)

Asbóth, Oroszlány, & Pályi Lecture Notes in Physics (2016)

- ∃ two sublattices ∃ a "canonical basis" where *H* is purely off-diag: $H = \begin{pmatrix} 0 & h^{\dagger} \\ h & 0 \end{pmatrix}$
- Chiral symmetry: $\Gamma H \Gamma = -H$ (Γ : unitary, Hermitian, local) $\Gamma = \sigma_z$
- In momentum space: $H_k = E_k \mathbf{n}_k \cdot \boldsymbol{\sigma}$ with $\mathbf{n}_k \perp \hat{\mathbf{z}}$ $\forall k$
- Winding:





- Bulk-edge correspondence: open-ended chains have $2\mathcal{W}$ edge modes

The winding W

• $\ensuremath{\mathcal{W}}$ may be calculated:

$$H_k = E_k \mathbf{n}_k \cdot \boldsymbol{\sigma}$$

• from n:
$$\mathcal{W} = \oint \frac{\mathrm{d}k}{2\pi} (\mathbf{n} \times \partial_k \mathbf{n})_z = \oint \frac{\mathrm{d}k}{2\pi} \partial_k \phi$$

• from the *eigenstates*:
$$\mathcal{W} = \oint \frac{\mathrm{d}k}{\pi} \mathcal{S}, \qquad \qquad \mathcal{S} = i \langle \psi_+ | \partial_k \psi_- \rangle$$

skew polarization

What if the Hamiltonian is not known?
 Can one *measure* the winding?

Yes, and it's simple!

Evolution in real time

Initial condition
 localized on the m=0 cell:



• Mean Chiral Displacement: $C(t) \equiv 2\langle \widehat{\Gamma m}(t) \rangle = 2 \left| \langle m_A(t) \rangle - \langle m_B(t) \rangle \right|$

$$\mathcal{C}(t) = \oint \frac{\mathrm{d}k}{2\pi} \langle \psi(t) | 2\sigma_z(i\partial_k) | \psi(t) \rangle = \mathcal{W} - \oint \frac{\mathrm{d}k}{2\pi} \cos(2\epsilon_k t) \partial_k \phi \xrightarrow{t \to \infty} \mathcal{W}$$

- Bulk measurement
- Fast convergence to $\ensuremath{\mathcal{W}}$
- Signals topological transitions!



Cardano, D'Errico, Dauphin, ... Marrucci, Lewenstein & PM Nature Comm. (2017)

Adding disorder to a topological insulator



disorder strength

Meier, An, Dauphin, Maffei, PM, Taylor and Gadway, Science (2018)

Atomic wires

• Atomic, non-interacting BEC

· Laser-driven coupling of discrete-momentum states

• 1D Hubbard model with built-in chiral symmetry:

$$H_{\text{eff}} \approx \sum_{j} t_j (e^{i\varphi_j} |\tilde{\psi}_{j+1}\rangle \langle \tilde{\psi}_j | + \text{h.c.})$$

• Full control on each tunneling's strength and phase







Detecting topology

A topological wire becomes trivial by adding disorder



disorder strength

color map: real-space computation of the winding

red line: critical boundary (diverging localization length)

Topological Anderson transition

A trivial wire is driven into the topological phase by adding disorder



disorder strength



Interacting fermionic chains



Half-filled Fermi sea $(U_{\perp}=U_{\parallel}=0)$

The MCD equals the winding straight away (no oscillations!)

[Haller, PM and Rizzi, on the arXiv very soon]

Interacting fermionic chains



Peierls-Hubbard model: U_{\perp} only ($U_{\parallel}=0$)



Interacting fermionic chains







Photonic quantum walks





1D

2D

Cardano, D'Errico, ..., and PM, Nature Comm. 2017 D'Errico, Di Colandrea, PM, ..., and Cardano, arXiv 2020

> see Francesco's poster this afternoon, and Alessio's talk on Thursday morning

quasi-periodicity of the energy spectrum leads to two inequivalent invariants \longleftrightarrow edge states

D'Errico, Cardano, Esposito, ..., PM et al., Optica 2020

see earlier talk by Filippo



Collaborators



Theory

Experiments



Maria Maffei

H Institut de Ciències Fotòniques



JOHANNES GUTENBERG UNIVERSITÄT MAINZ

Universität zu Köln

AARHUS UNIVERSITY



Andreas Haller

atomic wires







Fangzhao An

PHYSICS ILLINOIS Τ



Hughes Taylor







Alessio D'Errico

Filippo Cardano





Francesco di Colandrea

Chiara Esposito

SLAM group



Lorenzo Marrucci







Nathan Goldman



Arturo Camacho-Guardian

Georg Bruun

Matteo Rizzi



Eric J. Meier

Bryce Gadway



Conclusions

- The *mean chiral displacement* is a topological marker capturing the winding of chiral systems (static, periodically driven, disordered, and interacting)
- Experimental observation of a topological Anderson transition
- Detect topology & symmetry-breaking in interacting fermionic chains
- Characterization of a **periodically-driven chiral model** (two independent invariants \leftrightarrow two kinds of edge states)
- Dynamical observables for other topological classes?

Floquet chiral models:Nature CommMCD theory:New J. Phys.Topo. Anderson Insulator:Science 2018Topological polarons:Phys. Rev. BQuenched QWs:arXiv Jan 202Interacting chains:to be posted of

Nature Comm. 2017 New J. Phys. 2018 Science 2018 Phys. Rev. B (Rapid Comm.) 2019 arXiv Jan 2020 to be posted on arXiv very soon



Resistance to disorder



the MCD stays locked to the topological invariant as long as $\Delta{<}\Delta_{\rm gap}$

Higher windings

• Extension to long-ranged models:





 At critical boundaries: MCD converges to the mean of the winding in the neighboring phases

> Maffei, Dauphin, Cardano, Lewenstein & PM New J. Phys. 2018

Experimental proposal



- tilted optical lattice blocks tunneling
- inter-cell tunneling J' restored by two-photon transitions
- intra-cell tunneling implemented by RF transitions